

Relativity without light: A new proof of Ignatowski's theorem*

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Abstract: Ignatowski (1910) showed that assumptions about light are not needed to obtain Lorentzian kinematics as one of only few possibilities. We give a much simplified proof of his result as formulated by Gorini (1971) for $n+1$ -dimensional space-time.

* <http://arxiv.org/abs/2007.09301>

- Ignatowski (1910) was the first person to show that assumptions about light are not needed to establish special relativity as one of very few possibilities.
- His argument is a classic example of how *group theory* is powerful to describe Nature.
- New proof illustrates what can be done with *Lie theory* and a little *representation theory*.

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The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

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$$G = \left\{ \begin{pmatrix} A & 0 \\ 0 & \pm 1 \end{pmatrix} \underbrace{\exp \begin{pmatrix} 0 & b \\ \sigma \bar{b} & 0 \end{pmatrix}} : \begin{matrix} A \in O_3(\mathbf{R}) \\ b \in \mathbf{R}^3 \end{matrix} \right\} \quad \text{where} \quad \sigma = \frac{1}{c^2}$$

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$$\text{when} \quad \begin{aligned} b/c &= \vartheta e_1 \\ v/c &= \tanh \vartheta \end{aligned}$$

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When $\sigma = 0$ this still makes sense and is called the *Galilei* group.

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G first emerged in studies on the symmetry of Maxwell's equations:

Voigt (1887)

Larmor (1900)

Lorentz (1904)

Larmor: "For other physical quantities such as electric and magnetic forces, a less direct method must be followed, and will most probably be little by little and along the same lines. It is probable that the study of the invariance of the electromagnetic equations will give the key to the discovery of the laws governing the invariance of the other physical quantities."

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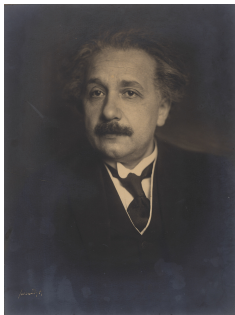
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Einstein (1905)



Poincaré (1905)



Minkowski (1908)

$G \cong O_{3,1}(\mathbf{R})$ — and a *group*.

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W. v. Ignatowsky (Berlin), Einige allgemeine Bemerkungen zum Relativitätsprinzip.

Als Einstein seinerzeit das Relativitätsprinzip einführte, nahm er parallel mit demselben an, daß die Lichtgeschwindigkeit c eine universelle Konstante sei, d.h. für alle Koordinatensysteme denselben Wert behalte. Auch Minkowski ging bei seinen Untersuchungen von der Invariante $r^2 - c^2 t^2$ aus, obwohl nach seinem Vortrage „Raum und Zeit“⁽¹⁾ zu urteilen, er dem c mehr die Bedeutung einer universellen Raum—Zeit-Konstante beilegte, als diejenige der Lichtgeschwindigkeit.

Nun habe ich mir die Frage gestellt, zu welchen Beziehungen bzw. Transformationsgleichungen man kommt, wenn man nur das Relativitätsprinzip an die Spitze der Untersuchung stellt und ob überhaupt die Lorentzschen Transformationsgleichungen die einzigen sind, die dem Relativitätsprinzip genügen.

“MINKOWSKI regarded c as, more than the speed of light, a universal space-time constant. Now I asked myself (...) whether the LORENTZ transformations are the only ones satisfying the principle of relativity.”

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group

Ignatowski's
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This is not just an academic question, for physicists *measure* whether light actually does travel at the invariant speed c (or in other words, whether the photon mass is exactly zero):

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Photon and graviton mass limits

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(Published 23 March 2010)

Efforts to place limits on deviations from canonical formulations of electromagnetism and gravity have probed length scales increasing dramatically over time. Historically, these studies have passed through three stages: (1) testing the power in the inverse-square laws of Newton and Coulomb, (2) seeking a nonzero value for the rest mass of photon or graviton, and (3) considering more degrees of

The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

The Lorentz
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Ignatowski's
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Gorini's
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Ignatowski assumes that “relativity” is implemented by a **group** of transformations

$$\begin{pmatrix} x' \\ t' \end{pmatrix} = \underbrace{\begin{pmatrix} a & b \\ c & d \end{pmatrix}}_g \begin{pmatrix} x \\ t \end{pmatrix}. \quad (1)$$

As (1) maps the line $x = 0$ to the line $x' = \frac{b}{d}t'$, he calls $v := \frac{b}{d}$ the *velocity* of the transformation. He requires that the velocity $\frac{-b}{a}$ of the inverse transformation $\frac{1}{ad-bc} \begin{pmatrix} d & -b \\ -c & a \end{pmatrix}$ be opposite to v , whence

$$a = d. \quad (2)$$

Now in order that a product $\begin{pmatrix} a & b \\ c & a \end{pmatrix} \begin{pmatrix} a' & b' \\ c' & a' \end{pmatrix} = \begin{pmatrix} aa'+bc' & * \\ * & aa'+b'c \end{pmatrix}$ still satisfy (2) we must have $bc' = b'c$, i.e. there is a constant σ such that

$$\frac{c}{b} = \frac{c'}{b'} = \sigma. \quad (3)$$

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So our transformations have the form $g = \begin{pmatrix} a & b \\ \sigma b & a \end{pmatrix}$ with $b = av$. Next Ignatowski argues that G should contain the reflection $R = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$ and hence also the product $RgRg = \begin{pmatrix} a^2 - \sigma b^2 & 0 \\ 0 & a^2 - \sigma b^2 \end{pmatrix}$, but no scalar matrices other than the identity (physics is not dilation-invariant). Hence we should have $1 = a^2 - \sigma b^2 \stackrel{b=av}{=} a^2(1 - \sigma v^2)$ and therefore

$$a = \frac{1}{\sqrt{1 - \sigma v^2}}. \quad (4)$$

Summing up, we have obtained

$$g = \begin{pmatrix} \frac{1}{\sqrt{1 - \sigma v^2}} & \frac{v}{\sqrt{1 - \sigma v^2}} \\ \frac{\sigma v}{\sqrt{1 - \sigma v^2}} & \frac{1}{\sqrt{1 - \sigma v^2}} \end{pmatrix} \quad (5)$$

which is the Lorentz transformation — or Galileo's, if $\sigma = 0$.

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which is the Lorentz transformation — or Galileo's, if $\sigma = 0$.

The Lorentz
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Ignatowski's
argument

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So our transformations have the form $g = \begin{pmatrix} a & b \\ \sigma b & a \end{pmatrix}$ with $b = av$. Next Ignatowski argues that G should contain the reflection $R = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$ and hence also the product $RgRg = \begin{pmatrix} a^2 - \sigma b^2 & 0 \\ 0 & a^2 - \sigma b^2 \end{pmatrix}$, but no scalar matrices other than the identity (physics is not dilation-invariant). Hence we should have $1 = a^2 - \sigma b^2 \underset{b=av}{=} a^2(1 - \sigma v^2)$ and therefore

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Grundlagen zu einer Theorie der Lorentztransformationen.

Von EMIL HAHN in Straßburg i. Els.

Einleitung.

Die Welt des von Einstein¹⁾ und Minkowski²⁾ eingeführten Raum-Zeitgebildes hat auf den ersten Blick etwas Befremdendes und, wie Minkowski³⁾ selbst bemerkt hat, Mißfälliges. Dieses Befremden schwindet sehr, wenn man den Voraussetzungen nachgeht, auf denen das Ganze beruht. Der entscheidende Schritt in dieser Richtung wurde von Ignatowsky⁴⁾ getan, der den Beweis zu führen suchte, daß das An-

The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

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Linear Kinematical Groups

VITTORIO GORINI★

Institut für Theoretische Physik (I) der Universität Marburg, Germany

Received November 12, 1970

Abstract. We prove a theorem which states that in an $(n+1)$ -dimensional space-time ($n \geq 3$) the only linear kinematical groups which are compatible with the isotropy of space are the Lorentz and Galilei groups. The special cases $n=1$ and $n=2$ are also briefly discussed.

Theorem (Gorini 1971)

Suppose $n \geq 2$ and let G be a subgroup of $GL_{n+1}(\mathbb{R})$ such that

$$G \cap \begin{pmatrix} GL_n(\mathbb{R}) & 0 \\ 0 & \mathbb{R}^\times \end{pmatrix} = \begin{pmatrix} O_n(\mathbb{R}) & 0 \\ 0 & \pm 1 \end{pmatrix}. \quad (6)$$

Write K for the right-hand side of (6). Then either $G = K$ or there is a number $\sigma \in \mathbb{R} \cup \{\infty\}$ such that

$$G = K \exp(p_\sigma), \quad (7)$$

where $p_\sigma = \left\{ \begin{pmatrix} 0 & b \\ \sigma \bar{b} & 0 \end{pmatrix} : b \in \mathbb{R}^n \right\}$ and $p_\infty = \left\{ \begin{pmatrix} 0 & 0 \\ c & 0 \end{pmatrix} : c \in \mathbb{R}^n \right\}$.

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In other words, G under the hypotheses must be isomorphic to one of these 5 possibilities — all of them well-known and named:

The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

- 1 if $\sigma > 0$, the Lorentz group $O_{n,1}(\mathbf{R})$ (Poincaré, 1906);
- 2 if $\sigma = 0$, the homogeneous Galilei group (Frank, 1908);
- 3 if $\sigma < 0$, the orthogonal group $O_{n+1}(\mathbf{R})$ (Jordan, 1870);
- 4 if $\sigma = \infty$, the homogeneous Carroll group (Lévy-Leblond, 1965);
- 5 if $G = K$, the homogeneous Aristotle group (Souriau, 1970).

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Gorini's result is remarkable, but didn't catch on — perhaps because his *proof* was neither concise nor enlightening.

The Lorentz
group

gnatowski's
argument

Gorini's
theorem

Our proof

156

V. Gorini:

By Lemma 1 (see (4.2)), we have

$$\frac{d(v) + v c(v) \cos \varepsilon}{d(v) + v c(v)} > 0, \quad \varepsilon \in [0, 2\pi), \quad (4.25)$$

and this gives the inequality

$$|v c(v)/d(v)| < 1, \quad v \in \mathcal{V}. \quad (4.26)$$

Defining

$$\cos \gamma(\varepsilon, v) = \frac{d\varepsilon(1 - \cos \varepsilon)}{[d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}}$$

and

$$\sin \gamma(\varepsilon, v) = \frac{a(d + v c) \sin \varepsilon}{[d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}},$$

with $\sin \gamma(0, v) = \sin \gamma(\pi, v) = 1$ and $\cos \gamma(0, v) = \cos \gamma(\pi, v) = 0$, we calculate

$$\begin{aligned} X(\varepsilon, v) &= C^{(1, 2)}(\pi - \gamma(\varepsilon, v)) L(a, v) C^{(1, 2)}(-\gamma(\varepsilon, v)) \\ &= \begin{pmatrix} \frac{d + v c \cdot \cos \varepsilon}{d + v c} & \varphi_1(a, c, d, e, v, \varepsilon) & 0 \dots 0 & \varphi_2(a, c, d, e, v, \varepsilon) \\ 0 & 1 & 0 \dots 0 & 0 \\ 0 & 0 & 1 \dots 0 & 0 \\ \hline 0 & 0 & 0 \dots 1 & 0 \\ \varphi_3(a, c, d, e, v, \varepsilon) & \varphi_4(a, c, d, e, v, \varepsilon) & 0 \dots 0 & \frac{d + v c \cdot \cos \varepsilon}{d + v c} \end{pmatrix}, \quad (4.27) \end{aligned}$$

where

$$\begin{aligned} \varphi_1(a, c, d, e, v, \varepsilon) &= \frac{\sin \varepsilon \{ d e^2 (1 - \cos \varepsilon) [d^2 \varepsilon^2 (1 - \cos \varepsilon) + a^2 (d + v c)^2 \cos \varepsilon] }{a e (d + v c) [d^2 \varepsilon^2 (1 - \cos \varepsilon)^2} \\ &\quad - a^2 (d + v c) [a^2 (d + v c)^2 \sin^2 \varepsilon - d e^2 (1 - \cos \varepsilon) (d \cdot \cos \varepsilon + v c)] \}, \quad (4.27a) \\ \varphi_2(a, c, d, e, v, \varepsilon) &= \frac{-v [d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}}{e (d + v c)}, \quad (4.27b) \end{aligned}$$

$$\varphi_3(a, c, d, e, v, \varepsilon) = \frac{c e [2d(1 - \cos \varepsilon) + v c \sin^2 \varepsilon]}{(d + v c) [d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}}, \quad (4.27c)$$

Linear Kinematical Groups

157

and

$$\varphi_4(a, c, d, e, v, \varepsilon) = \frac{c \sin \varepsilon (1 - \cos \varepsilon) [e^2 d - a^2 (d + v c)]}{a (d + v c) [d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}}. \quad (4.27d)$$

One checks that $d(X(\varepsilon, v)) = 1$, hence $X(\varepsilon, v)$ is an element of \mathcal{X} and, by Lemma 3, $X(\varepsilon, v) \in \mathcal{N}$. Therefore,

$$X(\varepsilon, v) = N(w(\varepsilon, v)), \quad (4.28)$$

where, by (4.27b),

$$w(\varepsilon, v) = \frac{v [d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon]^{1/2}}{e (d + v c \cos \varepsilon)}. \quad (4.29)$$

Eqs. (4.14) and (4.15) imply

$$\varphi_1(a, c, d, e, v, \varepsilon) = \varphi_4(a, c, d, e, v, \varepsilon) = 0. \quad (4.30)$$

One checks that for (4.30) to be satisfied it is necessary and sufficient that

$$d(\varepsilon) e^2 (v) = a^2 (v) [d(v) + v c(v)]. \quad (4.31)$$

Using (4.31) we get

$$d^2 \varepsilon^2 (1 - \cos \varepsilon)^2 + a^2 (d + v c)^2 \sin^2 \varepsilon = d e^2 [2d(1 - \cos \varepsilon) + v c \sin^2 \varepsilon]. \quad (4.32)$$

We have from (4.27) and (4.28), $\forall \varepsilon \in [0, 2\pi]$,

$$e(w(\varepsilon, v)) = f(w(\varepsilon, v)) = 1, \quad (4.33a)$$

$$a(w(\varepsilon, v)) = d(w(\varepsilon, v)) = \frac{d + v c \cos \varepsilon}{d + v c} \quad (4.33b)$$

and, using (4.32),

$$c(w(\varepsilon, v)) = \frac{c \{ d [2d(1 - \cos \varepsilon) + v c \sin^2 \varepsilon] \}^{1/2}}{d(d + v c)} \quad (4.33c)$$

and

$$w(\varepsilon, v) = \frac{v [d \{ 2d(1 - \cos \varepsilon) + v c \sin^2 \varepsilon \}]^{1/2}}{d + v c \cos \varepsilon}. \quad (4.33d)$$

From (4.33b, c, d) we get, $\forall \varepsilon \in [0, 2\pi]$,

$$1 + w^2(\varepsilon, v) \frac{c(v)}{v d(v)} > 0, \quad (4.34)$$

$$a(w(\varepsilon, v)) = \left(1 + \frac{c(v)}{v d(v)} w^2(\varepsilon, v) \right)^{-1/2} \quad (4.35)$$

and

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The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

Step 1: Lie group structure of G

G admits a canonical (“initial”) Lie group structure having Lie algebra

$$\mathfrak{g} := \{Z \in \mathfrak{gl}_{n+1}(\mathbf{R}) : e^{tZ} \in G \text{ for all } t \in \mathbf{R}\}. \quad (8)$$

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$$\text{Ad} \begin{pmatrix} R & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A & b \\ \bar{c} & d \end{pmatrix} = \begin{pmatrix} RAR^{-1} & Rb \\ \overline{Rc} & d \end{pmatrix}, \quad R \in O_n(\mathbf{R}). \quad (9)$$

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G admits a canonical (“initial”) Lie group structure having Lie algebra

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The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

Step 2: Determination of \mathfrak{g}

Either $\mathfrak{g} = \mathfrak{k}$ or there is $\sigma \in \mathbf{R} \cup \{\infty\}$ such that $\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{p}_\sigma$.

The Lorentz
group

Ignatowski's
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In fact we claim that all $Z = \begin{pmatrix} 0 & b \\ \bar{c} & 0 \end{pmatrix} \in \mathfrak{g} \cap M_3$ have b and c collinear.

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As this is contained in $[\mathfrak{g}, [\mathfrak{g}, \mathfrak{t}]] \subset \mathfrak{g}$, the lower right entry must be 0: so the Cauchy-Schwarz bound is attained, i.e. b and c are collinear.

So each $Z \in \mathfrak{g} \cap M_3$ is in some \mathfrak{p}_σ . Moreover σ must be the same for any two nonzero members Z_1, Z_2 : else, one would readily find — by considering linear combinations of Z_1 and $\text{Ad}(k)(Z_2)$ with $k \in K$ — that $\mathfrak{g} \cap M_3$ fills all M_3 , hence $\mathfrak{g} = \mathfrak{k} \oplus M_3$, which by (10) isn't a Lie subalgebra. So Step 2 is complete. \square

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Step 3: Passage from \mathfrak{g} to G

With $\mathfrak{g} = \mathfrak{k} \oplus \mathfrak{p}_\sigma$ as in Step 2, we have $K \exp(\mathfrak{p}_\sigma) \subset G \subset N(\mathfrak{g})$, where

$$N(\mathfrak{g}) := \{a \in GL_{n+1}(\mathbf{R}) : a\mathfrak{g}a^{-1} \subset \mathfrak{g}\}. \quad (11)$$

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Proof. The first inclusion holds by hypothesis and definition of \mathfrak{g} , the second because *any* subgroup G of any Lie group is contained in the normalizer (11) of its Lie algebra (since $e^{t\mathfrak{g}}g^{-1} = ge^{tZ}g^{-1}$). \square

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Proof. The first inclusion holds by hypothesis and definition of \mathfrak{g} , the second because *any* subgroup G of any Lie group is contained in the normalizer (11) of its Lie algebra (since $e^{t\mathfrak{g}}g^{-1} = ge^{tZ}g^{-1}$). \square

Step 3: Passage from \mathfrak{g} to G

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Step 4: Determination of $N(\mathfrak{g})$

$$N(\mathfrak{k} \oplus \mathfrak{p}_\sigma) = \begin{cases} \mathbf{R}^\times K \exp(\mathfrak{p}_\sigma) & \text{if } \sigma \in \mathbf{R}^\times, \\ \left(\begin{smallmatrix} \mathbf{R}^\times & 0 \\ 0 & \mathbf{R}^\times \end{smallmatrix} \right) K \exp(\mathfrak{p}_\sigma) & \text{if } \sigma \in \{0, \infty\} \text{ or } \mathfrak{p}_\sigma = \{0\}. \end{cases} \quad (12)$$

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Proof omitted. (A simple computation using Schur's lemma.) \square

The Lorentz
group

Ignatowski's
argument

Gorini's
theorem

Our proof

Steps 3 and 4, plus our hypothesis implying that $\begin{pmatrix} \lambda & 1 \\ 0 & \mu \end{pmatrix}$ is not in G unless it is in K (so $\lambda, \mu = \pm 1$), complete the theorem's proof. \square

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More details at <http://arxiv.org/abs/2007.09301>.